

V BRUNEL Workshop on Random Matrix Theory (Brunel University – West London)

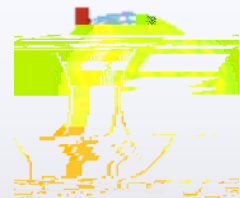


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by
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Thanks: R. Balbinot, S. Fagnocchi & I. Carusotto
(also for the figures)

- **Invariant Probability Distribution Function:**
- **Describe extended states (no localization)**
 - **Wigner statistics**
- **Gaussian Ensemble:**

Critical Statistics

Spontaneous Breaking of Invariance??

Weakly Confined Invariant Ensemble

$$\rho(E) \approx \frac{1}{\mathcal{N}(E)} \sum_{\alpha} \delta(E - E_{\alpha}) \quad \mathcal{N}(E) \propto \int_{-\infty}^{\infty} dx e^{-\kappa|x|} \delta(E - \lambda e^{\kappa|x|} \text{sign}(x))$$

- Non-Trivial density eigenvalue distribution
- Unfolding to make density constant:

$$\rho(E) \equiv \text{tr} \{ \delta(E - \mathbf{H}) \}$$

$$E_x = \lambda e^{\kappa|x|} \text{sign}(x)$$

$$\frac{1}{\mathcal{N}(E)} \delta(E - \mathbf{H}) = \frac{1}{\mathcal{N}(E_x)} \frac{dE_x}{dx} \delta(E - E_x)$$

Weakly Confined Invariant Ensemble

$$\langle Y_2(x, x') \rangle \approx \frac{1}{\mathcal{N}} \sum_{E \in \mathcal{E}} \frac{1}{\rho(E)} \left(\frac{1}{\rho(E)} \right) \approx \frac{1}{\mathcal{N}} \sum_{E \in \mathcal{E}} \frac{1}{\rho(E)^2}$$

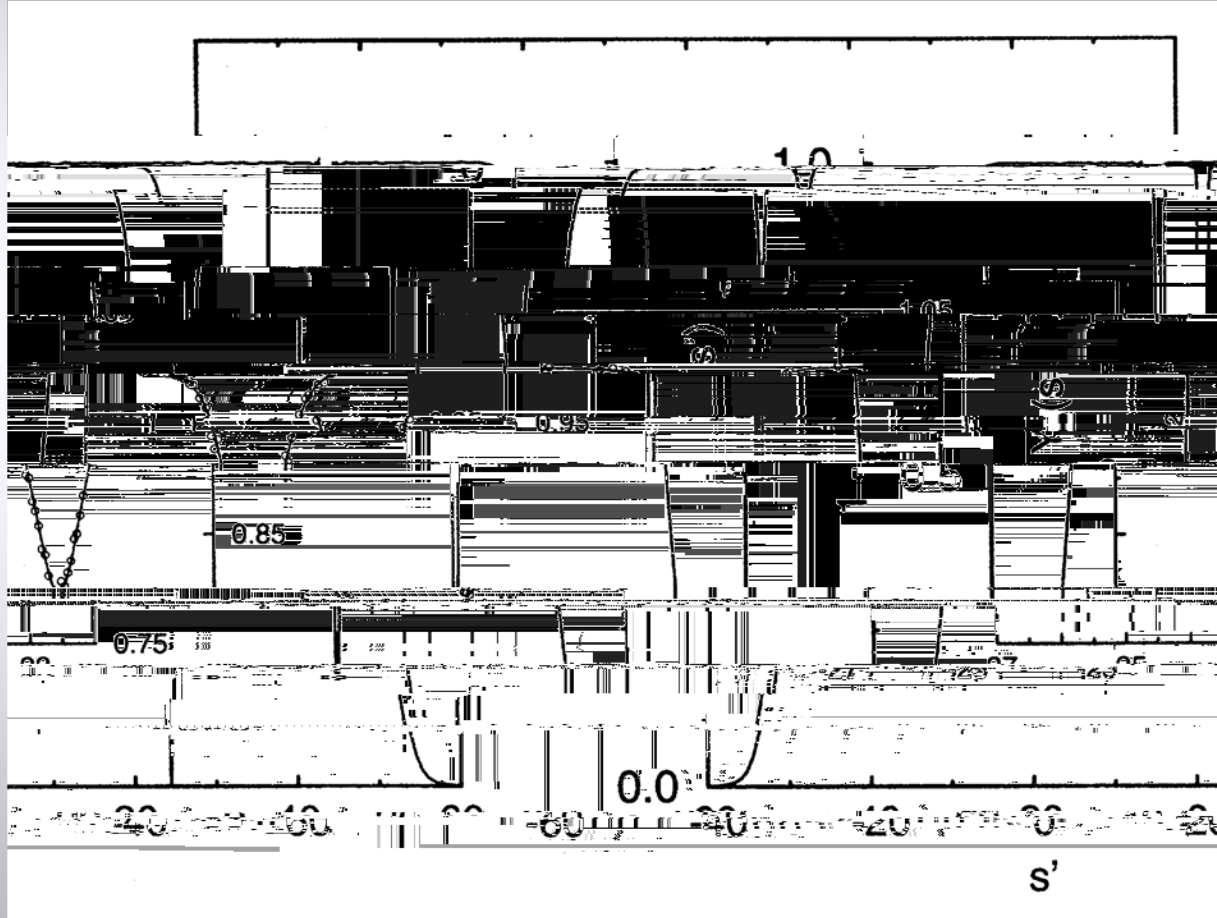
- For $e^{-2\pi^2/k} \ll 1$ semiclassical analysis (Canali et al '95):

$$\langle Y_2(x, x') \rangle \approx \delta(x - x') - \frac{\langle \rho(E_x) \rho(E_{x'}) \rangle - \langle \rho(E_x) \rangle \langle \rho(E_{x'}) \rangle}{\langle \rho(E_x) \rangle \langle \rho(E_{x'}) \rangle}$$

$$Y_2(x, x') \equiv \delta(x - x') - \frac{\langle \rho(E_x) \rho(E_{x'}) \rangle - \langle \rho(E_x) \rangle \langle \rho(E_{x'}) \rangle}{\langle \rho(E_x) \rangle \langle \rho(E_{x'}) \rangle}$$

Weakly Confined Invariant Ensemble

- Numerical check (Canali et al '95):



Acoustic Black-Hole

- Fluid pushed to move faster than it's speed of sound:



Hawking radiation

- Prediction: a Black Hole radiates particles with an **exact** thermal (*Black-Body*) spectrum
- Solid result due **only to horizon** (kinematical)
- Different ways to understand it:
 - Pair production close to horizon
 - Red-shifting of last escaping modes
 - Casimir effect
 - Bogoliubov overlap of positive frequency modes close to the horizon and at infinity
 - ...

QFT in Curved Space-Time

- Field quantization is basis-dependent:

$$\phi(x) = \sum_i [a_i f_i(x) + a_i^\dagger f_i^*(x)]$$

Plane waves

⇒ vacuum depends on the observer: $|0\rangle \neq |0'\rangle$

- For a different coordinate system:

$$\phi(x) = \sum_i [a_i f_i(x) + a_i^\dagger f_i^*(x)] = \sum_{i'} [\tilde{a}_{i'} \tilde{f}_{i'}(x) + \tilde{a}_{i'}^\dagger \tilde{f}_{i'}^*(x)]$$

⇒ $|0\rangle \neq |\tilde{0}\rangle$

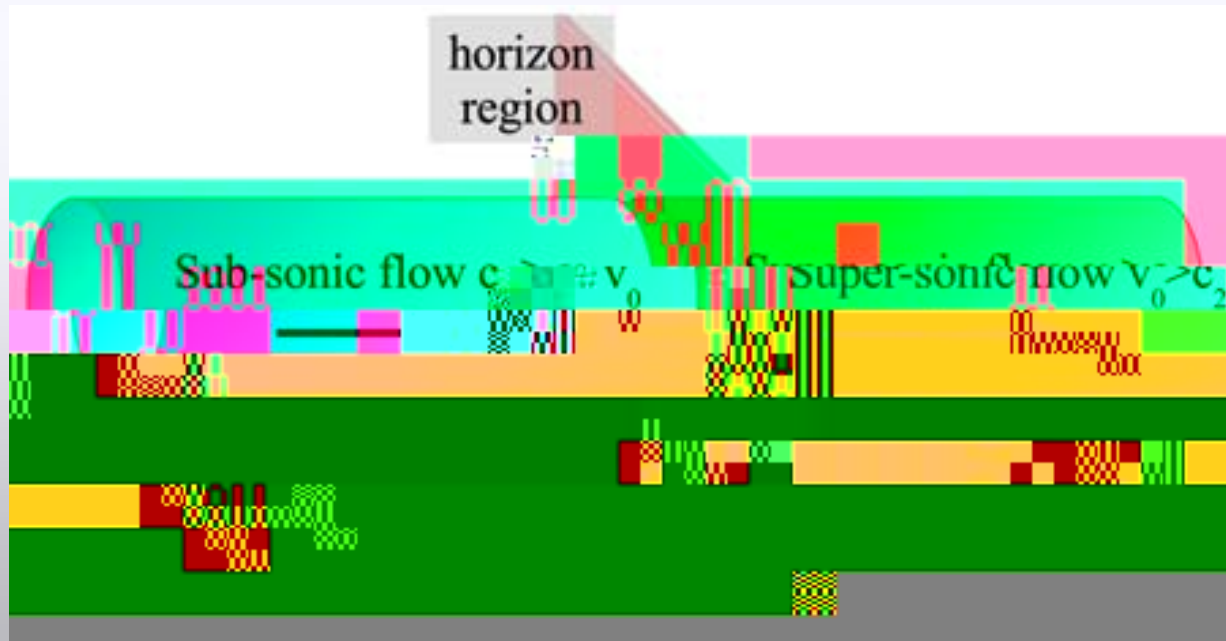
Hawking Radiation

- If



Acoustic BH in BEC

- Cool system \rightarrow Bose-Einstein Condensate
- Keep stream velocity constant & change speed of sound
- Effective dynamics is 1-D



Effective Gravity in fluids

- Low-Energy excitations (**phonons**) propagate on top of bulk stream
-

2-Point Correlator

- In flat space: $\langle \delta\phi(x, t) \delta\phi(x', t') \rangle \propto \ln(\Delta u^+ \Delta u^-)$

$$u^\pm \equiv t \pm \int \frac{dx}{c \mp v} \quad \leftarrow \text{Light-Cone coordinates}$$

and for the density: $\langle \delta\sigma(x) \delta\sigma(x') \rangle \propto \frac{1}{(x-x')^2}$

- Around the black hole: $u^- \rightarrow \tilde{u}^- = \frac{1}{\kappa} e^{-\kappa u^-} \text{erfc}(\kappa x)$

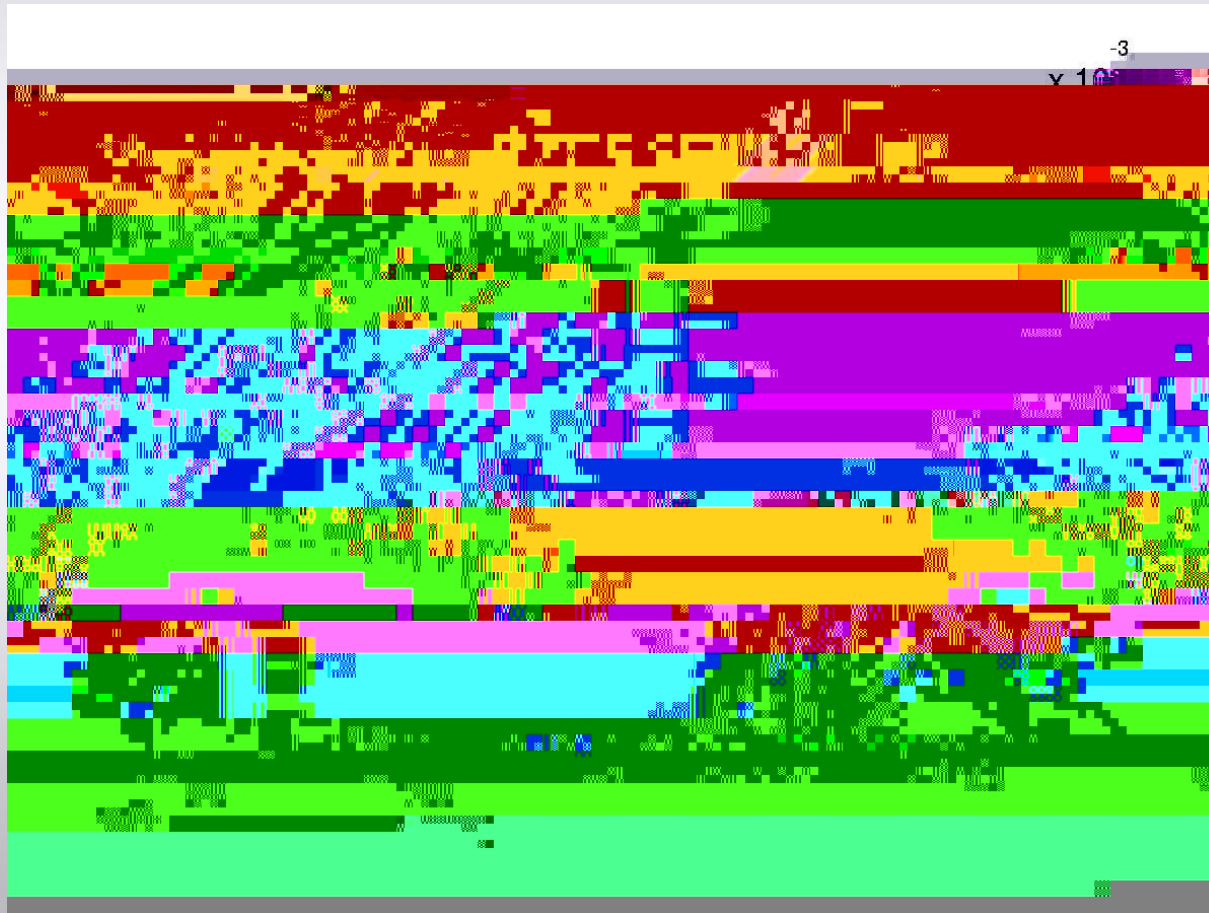
2-Point Correlator in curved metric

- Same non-local correlation as RME for $c_{r,l} = v \pm v/2$

(except for the oscillatory term...)

Numerical check

- Field theory prediction checked against ab-initio numerical simulation (Carusotto et al. '08)



- BEC system has non-local signature
- Low-Energy description in terms of free field in curved metric with horizon

Effective Theory for RME

$$\mathcal{L} = -\beta \sum_{n > m} \ln |E_n - E_m| + \sum_n V(E_n)$$

- Energy eigenvalues
 - coordinates of interacting particles
(fermions \Leftarrow level repulsion)
- Parametric evolution of RME
 - time coordinate
- Eigenvalue distribution
 - ground state configuration of 1D quantum model

Luttinger theory for RME

$$\langle \sigma_x \rangle = \frac{1}{\pi} \int_{-\pi}^{\pi} \frac{A_K}{2\pi \rho_0} \dots \rho(x, t) = \rho_0$$

- Two-Point function (Kravtsov et al. '00):

Unfolding:

$$\rho_0 = 1$$

$$V = \frac{1}{2} \int dx \Phi(x) \Phi(x')$$

- In flat space: $\langle \Phi(x, t) \Phi(x', t') \rangle \propto \ln(\Delta x^2 + \Delta t^2)$

$$\langle \Phi(x, t) \Phi(x', t') \rangle \propto \ln(\Delta x^2 + \Delta t^2)$$

2-Point Function
for Gaussian RME
(K=1: Unitary)

Non-invariant Critical Ensemble

- Critical Random Banded Matrix (Multifractal spectrum)

$$P(\mathbf{H}) \propto e^{-\sum_{n,m} A_{nm} |H_{nm}|^2} \quad A_{nm} = 1 + \frac{(n-m)^2}{2n}$$

- Thermal effective Luttinger Theory (Kravtsov & Tselik-2001)

g^* : critical conductance



$$\frac{1}{\Lambda^1}, T \equiv \frac{1}{g^*} Y_2(x, x') = T^2 \frac{\sin^2 \left[\frac{\pi}{2} (x - x') \right]}{\sin^2 \left[\frac{\pi}{2} (x + x') \right]}$$

Thermal Field Theory

- Diagonal part of 2-point function

$$T^2 \frac{\text{Tr} [\rho(x-x')]^2}{\sinh^2 [\pi T(x-x')]} Y_2(x, x') =$$

common to

- ∅ weakly confined invariant ensemble
 - ∅ Lorentzian banded matrix ensembles
 - ∅ Standard thermal field theory
- How to generate the non-translational invariant part?

Luttinger theory in curved metric

$$\mathcal{D}(\mathbf{H}) \approx \int \mathcal{D}\psi \mathcal{D}\bar{\psi} e^{-\text{Tr} V(\mathbf{H})} \quad V(E) \stackrel{|E| \rightarrow \infty}{\sim} \kappa \ln^2 |E|$$

- BEC system taught us that metric with horizons gives non-local correlation function

$$\mathcal{G}(x) \sim \frac{1}{\delta\pi\kappa J} \int_{-\infty}^{\infty} dt \sqrt{g(t)} \dots$$

- In 1+1 D any horizon metric can be approximated by

Rindler line element $g \equiv \det g^{\mu\nu}, ds^2 = g_{\mu\nu} d\xi^\mu d\xi^\nu$

$$ds^2 = -dt^2 + \frac{1}{y^2} dy^2 \leftarrow \text{Horizon at } y=0$$

- Let's choose: $y \equiv \sinh(\kappa x)$

Luttinger theory in Rindler space

- Far from the origin:

Luttinger Liquid in Rindler Space

- Remind two-Point function:

$$\begin{aligned}
 \langle \Phi(x) \Phi(x') \rangle &= \frac{1}{\pi^2} \left(\frac{\partial \Phi(x)}{\partial x} \frac{\partial \Phi(x')}{\partial x'} \right) \\
 + \dots &= -\frac{A_K^2}{2} \cos(2\pi(x-x')) \langle e^{i2\Phi(x)} e^{-i2\Phi(x')} \rangle.
 \end{aligned}$$

- With the new coordinates: $\left(\bar{x} = \frac{e^{\kappa|x|}}{2\kappa} \operatorname{sgn}(x) \right)$

$$\langle \Phi(x) \Phi(x') \rangle \propto \begin{cases} \ln \left[\frac{2}{\kappa} \sinh \frac{\kappa(x-x')}{2} \right], & x-x' > 0 \\ \ln \left[\frac{2}{\kappa} \cosh \frac{\kappa(x-x')}{2} \right], & x-x' < 0 \end{cases}$$

Luttinger Liquid in Rindler Space

- We recover exactly the RME correlation ($\mathbf{K}=1$):

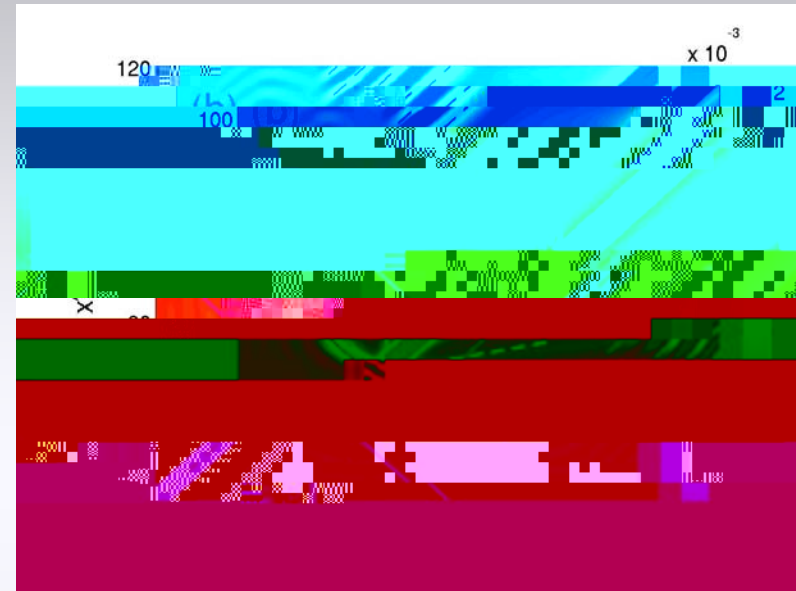
$$\langle \psi(x) \psi(x') \rangle \sim \frac{\kappa^2}{2\pi} \frac{\sin^2[\pi(x-x')]}{\cosh[\pi(x-x')]} \quad \text{for } \mathbf{K}=1$$

(Anomalous: non-translational invariant)

(Normal: translational invariant)

Summing up... (part 1)

- Luttinger Liquid predicts oscillatory term in correlator
- Possible to detect them in a BEC in Tonks-Girardeau regime



- We reproduced the **asymptotic 2-point function** in a Luttinger Liquid in curved space-time description
- Curved metric with horizons \rightarrow **Hawking radiation**
- Equivalence with **BEC system** (oscillatory term)
- Underlying integrable model as interesting probe for understanding Quantum Gravity (transformation problem)
- Strongly unrelaxed quantum condensate (topological fluctuations) → emergence of thermal bath